

Multipole Expansion Review

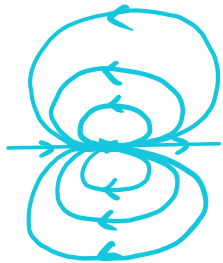
$$V(r) = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{r} + \frac{\vec{p} \cdot \hat{r}}{r^2} + \frac{Q_{ij} \hat{r}_i \hat{r}_j}{r^3} + \dots \right)$$

$\downarrow$ monopole $\downarrow$ dipole $\downarrow$ quadrupole

$$V_{\text{dipole}} = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2}$$

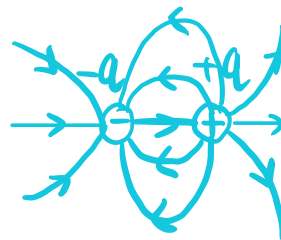
$$\begin{aligned}
 (\vec{E}_{\text{dipole}})_i &= -\partial_i V_{\text{dipole}} \\
 &= -\frac{1}{4\pi\epsilon_0} \partial_i \left[\frac{p_j \hat{r}_j}{r^2} \right] \\
 &= -\frac{p_j}{4\pi\epsilon_0} \partial_i \left(\frac{\hat{r}_j}{r^2} \right) \\
 &= -\frac{p_j}{4\pi\epsilon_0} \left[\frac{\partial_i \hat{r}_j}{r^2} + \hat{r}_j \partial_i \left(\frac{1}{r^2} \right) \right] \\
 &= -\frac{p_j}{4\pi\epsilon_0} \left[\frac{\delta_{ij}}{r^2} - 2\hat{r}_j \frac{1}{r^3} \partial_i r \right] \\
 &= -\frac{p_j}{4\pi\epsilon_0} \left[\frac{\delta_{ij}}{r^2} - \frac{3\hat{r}_i \hat{r}_j}{r^3} \right]
 \end{aligned}$$

$$\vec{E}_{\text{dipole}} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} [3(\vec{p} \cdot \hat{r})\hat{r} - \vec{p}]$$



"pure" dipole field

$$\nabla \cdot \vec{E} = 0$$



$$\vec{p} = q d \hat{e}$$

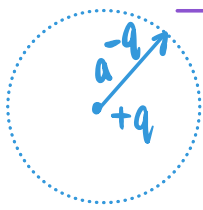
physical dipole

Insulator

Neutron atoms/molecules

intrinsic dipole moment?

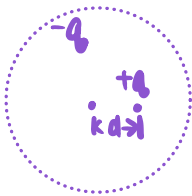
Polarization : Induce electric dipole
in response to external field



Without External Field:
Electrostatic System

Applying External Field:

Push positively charged electric cloud to the right
Push negatively charged electric cloud to the left



$$F_{ext} = E_{ext} q = \frac{Q^2}{4\pi\epsilon_0 a^3} d$$

$$Qd = 4\pi\epsilon_0 a^3 E_{ext}$$

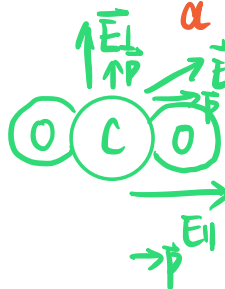
dipole
moment

$$\vec{P} = 4\pi\epsilon_0 a^3 \vec{E}_{ext}$$

$\alpha \leftarrow$ polarizability



$$\begin{aligned} \vec{E} &= \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{R}\right)^2 \hat{r} \\ &= \frac{Q}{4\pi\epsilon_0 R^2} \hat{r} \end{aligned}$$



$$P_i = \alpha_{ij} E_{ext,j}$$

α_{ij}
polarizability tensor

$$\vec{F} = \int d\tau \rho \vec{E}_{\text{ext}} = \vec{E}_{\text{ext}} \int d\tau \rho = 0$$

$$\begin{aligned} F_i &= \int d\tau \rho(\vec{r}) \vec{E}_{\text{ext}}(\vec{r}) = \int d\tau \rho(\vec{r}) \left[E_{\text{ext},i}(\vec{r}) + (\partial_j) \left[E_{\text{ext},i}(\vec{r}) + (\partial_j) E_{\text{ext},i}(\vec{r}) \right] \right. \\ &= \partial_t \left[\partial_i E_{\text{ext},i}(\vec{r}) \right] \underbrace{\int d\tau \rho(\vec{r})}_{\vec{P}_j} + \\ &= P_j \partial_j E_{\text{ext}}(\vec{r}) \end{aligned}$$

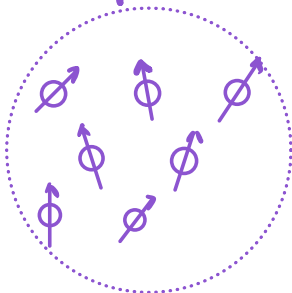
$$\vec{F} = (\vec{P} \cdot \nabla) \vec{E}_{\text{ext}} \quad \vec{F}_i = P_j \partial_j E_{\text{ext},i}$$

$$\begin{aligned} \vec{N} &= \int d\tau \vec{r} \times (\rho \vec{E}_{\text{ext}}) \\ &= \underbrace{\int d\tau \rho \vec{r}}_{\vec{P}} \times \vec{E}_{\text{ext}} \\ &= \vec{P} \times \vec{E}_{\text{ext}} \end{aligned}$$

Macroscopic field

vs.

Microscopic field



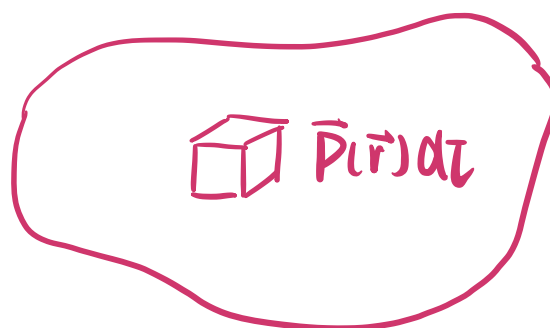
$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int d\tau' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|^2} (\vec{r} - \vec{r}')$$



$$\vec{E} = \vec{E}_{ave}$$

$$\vec{E}_{ave}(\vec{r}) = \frac{1}{\frac{4\pi}{3}R^3} \int d\tau' \vec{E}_{macro}(\vec{r}+\vec{r}')$$

$$\vec{E}(\vec{r}) = \sum_{i=1}^N \frac{1}{4\pi\epsilon_0} \frac{1}{|\vec{r}-\vec{r}_i|^3} \left[\frac{3\vec{p}_i(\vec{r}-\vec{r}_i)(\vec{r}-\vec{r}_i)}{|\vec{r}-\vec{r}_i|^2} - \vec{p}_i \right] \quad \leftarrow \text{Macroscopic field}$$



Polarization

$\vec{P} = \frac{\text{total molecular dipole momentum}}{\text{Volume}}$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_V d\tau' \frac{\vec{P}(\vec{r}') \cdot (\vec{r}-\vec{r}')}{|\vec{r}-\vec{r}'|^3} \quad \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} = -\nabla' \frac{1}{|\vec{r}-\vec{r}'|} = +\nabla' \frac{1}{|\vec{r}-\vec{r}'|}$$

$$= \frac{1}{4\pi\epsilon_0} \int_V d\tau' \vec{P}(\vec{r}') \cdot \nabla' \frac{1}{|\vec{r}-\vec{r}'|}$$

$$= \frac{1}{4\pi\epsilon_0} \int_V d\tau' \left[\nabla' \cdot \left(\frac{\vec{P}(\vec{r}')}{|\vec{r}-\vec{r}'|} \right) - \frac{1}{|\vec{r}-\vec{r}'|} \nabla' \cdot \vec{P}(\vec{r}') \right]$$

$$= \frac{1}{4\pi\epsilon_0} \oint_S d\vec{a}' \cdot \frac{\vec{P}(\vec{r}') \hat{n}'}{|\vec{r}-\vec{r}'|} - \frac{1}{4\pi\epsilon_0} \int_V d\tau' \frac{\nabla' \cdot \vec{P}(\vec{r}')}{|\vec{r}-\vec{r}'|}$$

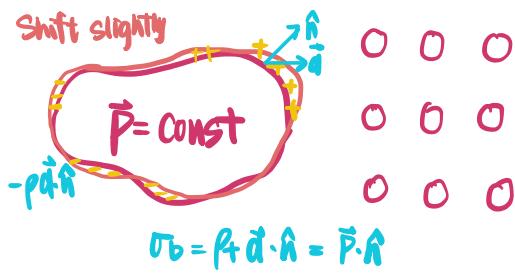
bound surface
charge

$$\sigma_b \equiv \vec{P} \cdot \hat{n}$$

$$= \frac{1}{4\pi\epsilon_0} \oint_S d\vec{a}' \frac{\sigma_b(\vec{r}')}{|\vec{r}-\vec{r}'|} + \frac{1}{4\pi\epsilon_0} \int_V d\tau' \frac{\rho_b(\vec{r}')}{|\vec{r}-\vec{r}'|}$$

bound volume
charge

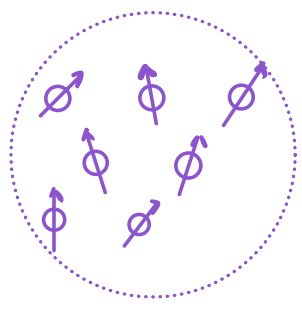
$$\rho_b = -\nabla \cdot \vec{P}$$



Same amount of separation

$$\rho_+ + \rho_- = 0$$

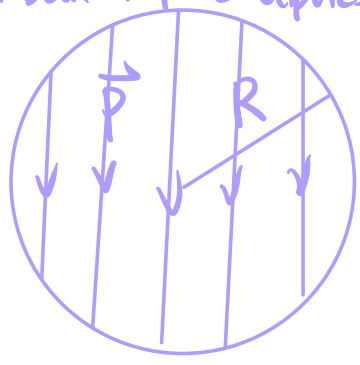
Things only happen on boundary



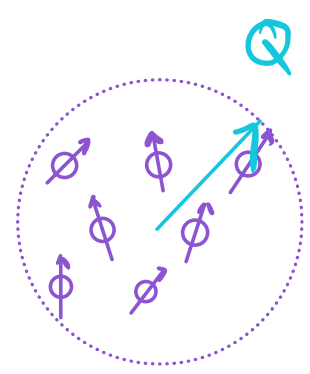
$$\vec{E} = \vec{E}_{\text{out}} + \vec{E}_{\text{in}}$$

$$\vec{E}_{\text{in}} = - \frac{\vec{P}}{3\epsilon_0} = - \frac{1}{3\epsilon_0} \frac{3}{4\pi R^3} \sum_{i=1}^N \vec{P}_i$$

uniform ball of pure dipoles



$$\vec{E}_{\text{in}} = - \frac{\vec{P}}{3\epsilon_0}$$



$$Q \vec{E}_{in} = \int d\vec{l} \rho(\vec{r}) \frac{Q}{4\pi\epsilon_0} \frac{\vec{r}}{R^3}$$

$$= \frac{Q}{4\pi\epsilon_0 R^3} \underbrace{\int d\vec{l} \rho(\vec{r}) \vec{r}}_{\Sigma P_i = \vec{P} \frac{4\pi}{3} R^3} = \frac{Q}{4\pi\epsilon_0 R^3} \frac{4\pi}{3} R^3 \vec{P}$$

$$\vec{E}_{in} = -\frac{\vec{P}}{3\epsilon_0}$$