

Isospin Symmetry

233A LEC2

$m_p \approx m_n$

$m_p \sim 1 \text{ GeV}$

$\Delta(m_n, m_p) \sim \text{MeV}$

• proton $\sim uud$

neutron $\sim udd$

(strong force) (symmetry)
QCD $\checkmark$ treats u & d same
EM - - - different

$$|N\rangle = \begin{pmatrix} |p\rangle \\ |n\rangle \end{pmatrix}$$

$$|N\rangle \rightarrow U|N\rangle$$

$$\begin{pmatrix} |p\rangle \\ |n\rangle \end{pmatrix} \rightarrow \begin{pmatrix} U_{11} & U_{12} \\ U_{21} & U_{22} \end{pmatrix} \begin{pmatrix} |p\rangle \\ |n\rangle \end{pmatrix}$$

$$U^\dagger U = \mathbb{I}$$

$$U \in \text{SU}(2)$$

$$U^\dagger U = \mathbb{I}$$

$$\det U = 1$$

Aside

$$\begin{pmatrix} |e, \uparrow\rangle \\ |e, \downarrow\rangle \end{pmatrix} \rightarrow U_{\text{space-time}} \times \begin{pmatrix} |e, \uparrow\rangle \\ |e, \downarrow\rangle \end{pmatrix}$$

Invariance

$$U^\dagger H U = H \quad \leftarrow \text{Hamiltonian}$$

Show that $m_n = m_p$

$$m_p \equiv \langle p | H | p \rangle$$

$$= \langle p | U^\dagger \underbrace{U H U^\dagger}_H U | p \rangle$$

$$\rightarrow m_p = \langle n | H | n \rangle = m_n$$

$$U = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \det U = +1$$

$$U | p \rangle = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} = | n \rangle$$

$SU(2)$

$$U_{\text{isospin}} = \mathbb{I} + i \sum_{a=1}^3 \delta\phi_a T^a$$

$\uparrow$ generators of $SU(2)$

$$T^a = \frac{1}{2} \tau^a$$

$\uparrow$ Pauli Matrices

$$(i) U^\dagger U = \mathbb{I}$$

$$\Rightarrow i \sum_{\alpha=1}^3 \delta\phi_\alpha [T^{\alpha\dagger} - T^\alpha] = 0$$

$$\Rightarrow T^{\alpha\dagger} = T^\alpha \rightarrow \text{Hermitian}$$

$$(ii) \det U_{\text{isospin}} = 1$$

$$U_{\text{isospin}} = e^{i\delta\phi_\alpha T^\alpha}$$

$$\det e^A = e^{\text{Tr} A}$$

$$\det U_{\text{isospin}} = e^{i\delta\phi_\alpha \text{Tr}[T^\alpha]} = 1$$

$$\text{Tr} T^a = 0 \rightarrow \text{Traceless}$$

$\Rightarrow$ Three independent, traceless Hermitian.

$$\tau_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \tau_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \tau_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

States

Eigenvalue

$$\begin{cases} |\vec{T}|^2 = T(T+1) \\ T_3 \end{cases}$$

	I	I_3
$ p\rangle$	$\frac{1}{2}$	$\frac{1}{2}$
$ n\rangle$	$\frac{1}{2}$	$-\frac{1}{2}$

$$T_3|n\rangle = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = -\frac{1}{2} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = -\frac{1}{2}|n\rangle$$

Ex:

	I	I_3
pp	1	1
$\frac{1}{\sqrt{2}}(pn+np)$	1	0
nn	1	-1
$\frac{1}{\sqrt{2}}(pn-np)$	0	0

$$|spin \frac{1}{2}\rangle \otimes |spin \frac{1}{2}\rangle$$

$$= |spin 1\rangle \otimes |spin 0\rangle$$

Pions for Isotriplets

$$T_i |\pi_j\rangle = i \epsilon_{ijk} |\pi_k\rangle$$

$$i=1,2,3$$

Define:

$$|\pi^\pm\rangle = \frac{1}{\sqrt{2}} (|\pi_1\rangle \pm i |\pi_2\rangle)$$

$$T_3 |\pi^\pm\rangle = \frac{1}{\sqrt{2}} [\epsilon_{312} |\pi_2\rangle \\ \pm (-1) \epsilon_{321} |\pi_1\rangle]$$

$$\begin{aligned} \epsilon_{312} &= 1 \\ \epsilon_{321} &= -1 \end{aligned} \quad = \frac{1}{\sqrt{2}} [|\pi_2\rangle \pm |\pi_1\rangle]$$
$$= \pm |\pi^\pm\rangle$$

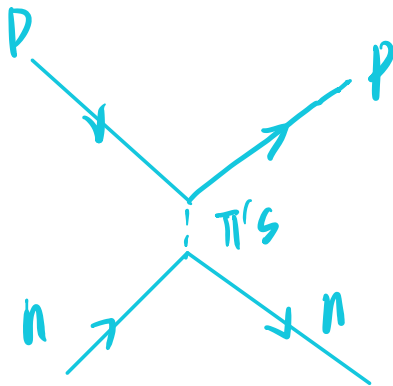
Eigenvalue ± 1

under the action of generator T_3

$$T_3 |\pi_2\rangle = i \epsilon_{33K} |\pi_K\rangle$$

$$= 0$$

	I	I ₃
π^+	1	+1
π^0	1	0
π^-	1	-1



Effective Lagrangian

$$\mathcal{L} \propto \bar{N}_i (T^a)_{ik} \pi^a N_k$$

$$\begin{array}{ll} \bar{i}=1,2 & N_1=p \\ a=1,2,3 & N_2=n \end{array}$$

$2 \oplus 2 = 3 \oplus 1$ Pion can transform like a fundamental

$$N \rightarrow UN \quad \bar{N} \rightarrow \bar{N}U^\dagger$$

$$\vec{T} \cdot \vec{\pi} \rightarrow U \vec{T} U^\dagger \vec{\pi}$$

$SU(2)$ Antifundamental
= fundamental

$$L \rightarrow \bar{N} U^\dagger U \vec{T} U^\dagger \vec{\pi} U N$$

$$= \bar{N} \vec{T} N = L$$

Lagrangian goes to itself under $SU(2)$

$$T^a \pi^a = \begin{pmatrix} \pi^3 & \sqrt{2} & \pi^- \\ \sqrt{2} & \pi^+ & -\pi^3 \end{pmatrix}$$

$$L \sim (\bar{p} \ \bar{n}) \begin{pmatrix} \pi^3 & \sqrt{2} & \pi^- \\ \sqrt{2} & \pi^+ & -\pi^3 \end{pmatrix} \begin{pmatrix} p \\ n \end{pmatrix}$$

$$\sim \sqrt{2} \bar{p} n \pi^+ + \sqrt{2} \bar{n} p \pi^- + (\bar{p} p - \bar{n} n) \pi^0$$

$$\pi^+ \rightarrow e^+ \bar{\nu}_e \pi^+ \quad \pi^- \rightarrow e^- \bar{\nu}_e \pi^-$$

$$\bar{p} \rightarrow \bar{p} e^- \bar{\nu}_e$$

Strangeness

"V-particles"

• K-mesons, Λ -hyperons

① $\pi^- + p \rightarrow K^0 + \Lambda$ Big cross section

$$\sigma(\pi^- p \rightarrow K^0 \Lambda) \approx 1 \text{ mb} = 10^{-27} \text{ cm}^2$$

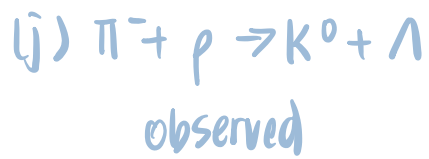
$$\sigma(\pi^- p \rightarrow \text{anything}) \approx 40 \text{ mb}$$

$$E_\pi \approx 1.5 \text{ GeV}$$

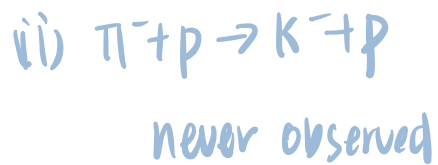
$$\sigma = (f \hbar)^2 \sim (10^{-15} \text{ cm})^2 = 10^{-26} \text{ cm}^2$$

$$f \hbar \sim \frac{1}{200 \text{ MeV}} \sim \frac{1}{m_\pi}$$

$$\hbar c = 1 = 200 \text{ MeV fm}$$



follow conservation of strangeness



update strangeness

Puzzle 2



Way too slow

SU(3)