

$$\begin{aligned}
 F(q, p, t) &= q \cdot p + F_2(q, p, t) \\
 &= q \cdot p + \varepsilon G(q, p, t) \\
 &= q \cdot p + \varepsilon G(q, p, t) + o(\varepsilon)
 \end{aligned}$$

$$(q, p) \rightarrow (Q, P)$$

Examples =

1) $G = H$ is the generator of infinitesimal time translations
 hence if H is time-independent then H is conserved

2) $G = p_i$ is the generator of infinitesimal translations of q_i
 hence q_i is constant of motion if H is invariant under translations of q_i

$$\begin{aligned}
 3) \quad G &= x p_y - y p_x \\
 &= (x p) \cdot \hat{z} = L_z
 \end{aligned}$$

is the generator of infinitesimal rotations about z-axis

L_z is conserved if H is invariant

under rotations around the
z-axis

If L_x, L_y are both constants
of motions then $\{L_x, L_y\} = L_z$

so L_z is also constant of motion
and so L is constant

$$\{L_i, L_j\} = \epsilon_{ijk} L_k$$

where ϵ_{ijk} are the structure
constant of the Lie algebra

$SO(3)$ associated

Recall that under the parametrized
family of ICIS with generating

function $G(z, \alpha)$

$$\delta z = \delta \alpha \int \frac{\partial G}{\partial z}$$

and so far any phase function $W(z|\alpha)$

$$\delta W = d\alpha \left(\frac{\partial W}{\partial \underline{z}} \right)^T \cdot \delta \underline{z}$$

$$= d\alpha \left(\frac{\partial W}{\partial \underline{z}} \right)^T \underline{J} \frac{\partial G}{\partial \underline{z}}$$

$$= d\alpha \{W, G\}$$

$$\frac{dW}{d\alpha} = \{W, G\} = -L_G W$$

where $L_G(\cdot) = \{G, (\cdot)\}$ Lie derivative

The properties of L_G follow the

properties of Poisson bracket

$$\begin{array}{l} \text{Linearity} \\ L_{\lambda A + \mu B} = \lambda L_A + \mu L_B \\ L_A(\lambda B + \mu C) = \lambda L_A B + \mu L_A C \end{array} \left. \vphantom{\begin{array}{l} \text{Linearity} \\ L_{\lambda A + \mu B} = \lambda L_A + \mu L_B \\ L_A(\lambda B + \mu C) = \lambda L_A B + \mu L_A C \end{array}} \right\} \lambda, \mu \text{ are constant}$$

Anti-symmetry

$$\text{Jacobi identity} \quad [L_A, L_B] = L_A L_B - L_B L_A = L_{\{A, B\}}$$

It follows that

$$W(z, a) = e^{\underbrace{\int_0^a da L_G(a)}_T} W(z, 0)$$

T is the evolution operator that evolves $W(z, 0) = W(z, a)$

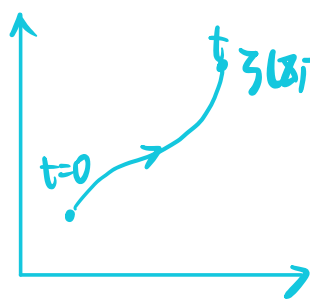
Properties: If $G=H$

T generates time translations

$$T^{-1}(t) = T(-t)$$

$$T^{-1}(t) T(t) = \text{Id}$$

$$T(t_1) T(t_0) = T(t_1 + t_0) \quad \text{Semigroup}$$



$$z(z, t) = Tz \quad \text{similarly } T W(z) = W(z(z))$$

$$\text{For example, } K(z(z)) = H(z)$$

$$\text{becomes } TK = H$$

$$K = T^{-1}H$$

Important remark:

$$e^{-\int_0^t dt' LH(t')} = 1 - \int_0^t dt' LH(t') \\ + \int_0^t dt' \int_0^{t'} dt'' LH(t'') LH(t') \dots$$

$$\int_0^t f(t') dt' \int_0^{t'} g(t'') dt'' + \int_0^t g(t') dt' \int_0^{t'} f(t'') dt'' \\ = \frac{1}{2} \left(\int_0^t f(t') dt' \right) \left(\int_0^t g(t'') dt'' \right)$$

This provides a generalization of Taylor series

$$f(x+a) = \left(\exp a \frac{d}{dx} \right) f(x) \\ = f(x) + a f'(x) + \frac{1}{2} a^2 f''(x) + \dots$$

Ex Suppose H is autonomous (time-independent)

$$e^{-tLH} = 1 - tLH + \frac{1}{2} t^2 LH^2 \\ = \text{id} - t \{H, \cdot\} + \frac{1}{2} t^2 \{H, \{H, \cdot\}\} -$$

Suppose $H = \frac{p^2}{2m} - \max$

particle of mass under constant acceleration
 a

$$x(t) = \left[id - t \{H, \cdot\} + \frac{1}{2} t^2 \{H, \{H, \cdot\}\} - \dots \right] x |_{t=0}$$

$$\{H, x\} = \frac{\partial H}{\partial x} \frac{\partial x}{\partial p} - \frac{\partial H}{\partial p} \frac{\partial x}{\partial x}$$

$$= -\frac{\partial H}{\partial p} = -\frac{p}{m}$$

$$\{H, x\} = -\frac{p}{m}$$

$$\{H, p\} = \frac{\partial H}{\partial x} \frac{\partial p}{\partial p} - \frac{\partial H}{\partial p} \frac{\partial p}{\partial x}$$

$$p = -m\{H, x\}$$

$$= \frac{\partial H}{\partial x} = -ma$$

$$\{H, \{H, x\}\} = -\frac{1}{m}(-ma) = a$$

$$x(t) = \left[id - t \{H, \cdot\} + \frac{1}{2} t^2 \{H, \{H, \cdot\}\} - \dots \right] x |_{t=0}$$

$$= x_0 + \frac{p_0}{m} t + \frac{a}{2} t^2$$

$$\text{where } (x_0, p_0) = (x(t), p(t)) |_{t=0}$$

Extended phase space

Suppose n -dof Hamiltonian $H(z,t)$ is non-autonomous

Variational Principle

$$\begin{aligned} \delta \int z dt &= \delta \int (p \cdot \dot{q} - H) dt \\ &= \delta \int \left(p \cdot \frac{dq}{dz} - H \frac{dt}{dz} \right) dz \quad \downarrow \text{time like variable} \\ &= \delta \int p \cdot \frac{dq}{dz} dz \end{aligned}$$

this suggest we define i

$$P_i = p_i \quad Q_i = q_i \quad i=1, \dots, n$$

$$Q_{n+1} = t$$

$$P_{n+1} = -H$$

$$i = n+1$$

This change of variation is guaranteed

by

$$\bar{F}_2 = \sum_i^n P_i q_i + P_{n+1} t$$

$$\text{s.t.} \quad p_i = \frac{\partial \bar{F}_2}{\partial q_i} \quad q_i = \frac{\partial \bar{F}_2}{\partial p_i}$$

$$K(Q, P) = H(q, p, t) + \frac{\partial \bar{F}_2}{\partial t} = H(q, p, t) - H$$